

Using the Magnetospheric Multiscale Mission to Examine Electric Fields in Turbulent Plasmas

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Key Points

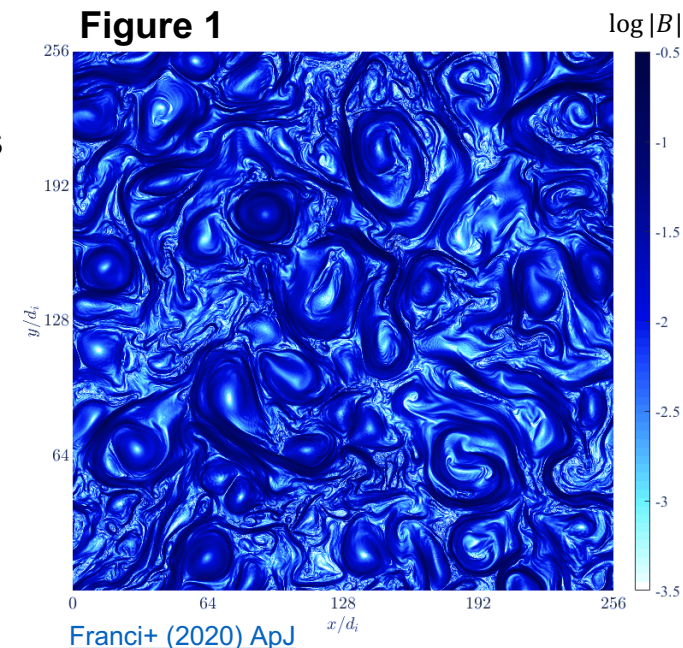
- Using NASA's Magnetospheric Multiscale, we examine in detail how generalised Ohm's Law shapes the turbulent electric field for the first time
- Results provide insight into the interplay between the Hall and electron pressure terms, which is important for understanding turbulent dissipation
- Ohm's Law allows a direct examination of the relative importance of linear and nonlinear dynamics

Turbulence is a fundamental process for particle energization in plasmas throughout the Universe, from the solar wind and planetary magnetospheres to accretion discs and galaxy clusters

Turbulence is characterised by **highly-nonlinear fluctuations across a wide range of length scales** (Fig 1)

Many space plasmas are collisionless, resulting in a variety of possible mechanisms for dissipating the fluctuation energy

Disentangling these dissipation mechanisms is a major open problem in plasma turbulence research



Analyzing Electric Fields in Collisionless Plasmas

Figure 2

$$\mathbf{E}_{Ohm} = \underbrace{-\mathbf{u} \times \mathbf{B}}_{\mathbf{E}_{MHD}} + \underbrace{\frac{1}{en} \mathbf{J} \times \mathbf{B}}_{\mathbf{E}_{Hall}} - \underbrace{\frac{1}{en} \nabla \cdot \mathbf{P}_e}_{\mathbf{E}_{Pe}} + \underbrace{\frac{m_e}{e^2 n} \nabla \cdot \left(\mathbf{u} \mathbf{J} + \mathbf{J} \mathbf{u} - \frac{\mathbf{J} \mathbf{J}}{en} \right) + \frac{m_e}{e^2 n} \frac{\partial \mathbf{J}}{\partial t}}_{\mathbf{E}_{Inertia}} + \underbrace{\sum_{\ell=1}^{\infty} \left(-\frac{m_e}{m_i} \right)^\ell \mathcal{M}_\ell}_{\mathbf{E}_{\delta m_e}}$$

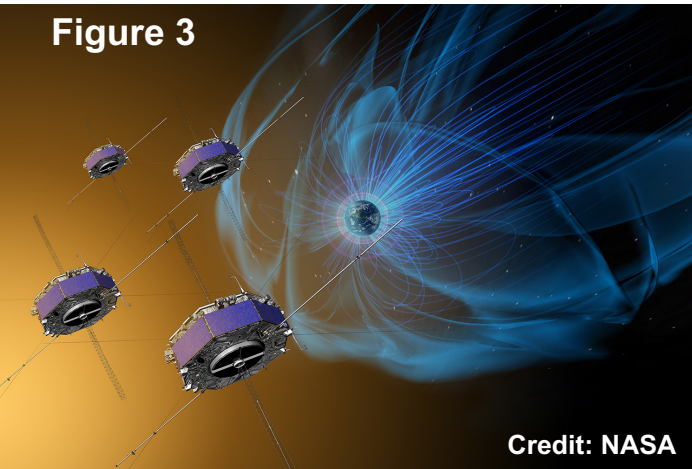
Magnetic Field "Frozen-In" to ion fluid
 Magnetic Field "Frozen-In" to electron fluid
 Correction due to electron thermal motions
 Correction due to difference between ion and electron inertia
 Finite electron mass corrections

$$\mathcal{M}_\ell = \frac{2}{en} \mathbf{J} \times \mathbf{B} - \frac{1}{en} \nabla \cdot (\mathbf{P}_e + \mathbf{P}_i) + \frac{m_e}{e^2 n} \left[\nabla \cdot \left(\mathbf{u} \mathbf{J} + \mathbf{J} \mathbf{u} - (1 + 2\ell) \frac{\mathbf{J} \mathbf{J}}{en} \right) + \frac{\partial \mathbf{J}}{\partial t} \right]$$

Since magnetic fields (\mathbf{B}) do no work, electric fields (\mathbf{E}) are required to exchange energy between \mathbf{B} and the particles $\rightarrow \mathbf{E}$ is key to understanding both turbulent dissipation and the nonlinear dynamics

In collisionless plasmas, \mathbf{E} is governed by a **generalised Ohm's Law** in which the different terms correspond to different dynamical processes (**Fig. 2**)

Figure 3



NASA's **Magnetospheric Multiscale (MMS)** mission consists of 4 closely spaced satellites, providing 3D, high-time-resolution, multipoint plasma measurements (**Fig. 3**)

MMS is uniquely capable of directly probing nearly all the terms in Ohm's Law down to length scales approaching those of electron motions [e.g. [Torbert+ \(2016\) GRL](#)]

How do the terms in Ohm's Law behave as a function of scale?

Three intervals of MMS data from Earth's magnetosheath are analysed
 [Results from one interval are shown here as an example]

Timeseries of each Ohm's Law term are computed and then the power spectra are compared to the power spectrum of the directly measured E (Fig. 4)

Frequencies are converted to wavenumbers using the average flow velocity $\rightarrow k = 2\pi f / U_0$

Results

Good agreement between measured E and E_{Ohm}

E_{MHD} provides dominant contribution at large scales, indicating B is frozen-in to ion fluid flows

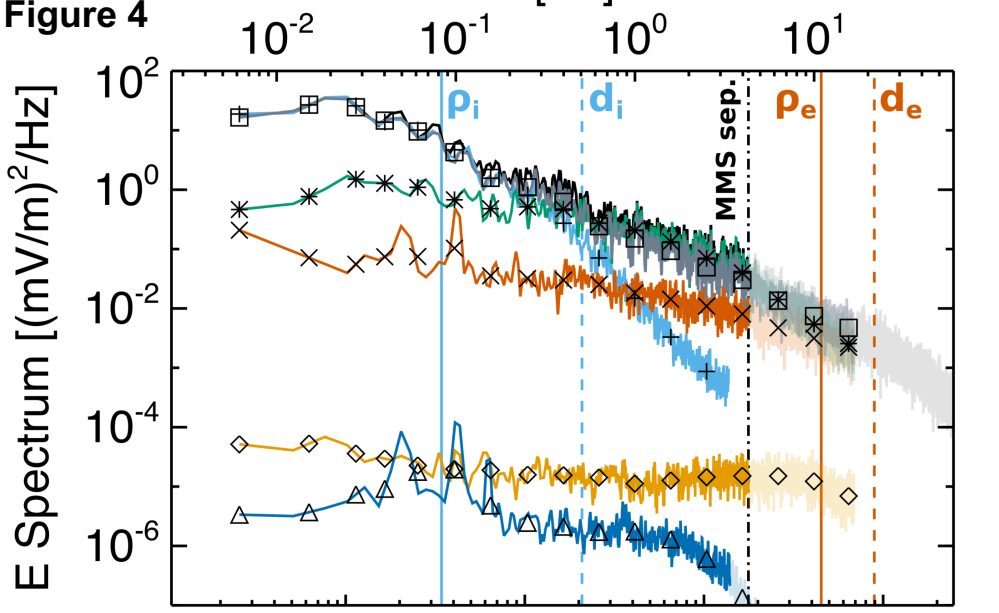
E_{Hall} makes largest contribution at sub-ion scales
 \rightarrow scale of transition typically occurs near

$$kd_i \sim \frac{\delta u_{rms}}{\delta b_{rms} / \sqrt{\mu_0 m_i n}}$$

E_{Pe} provides non-zero contribution to sub-ion scale electric field

$E_{inertia}$ and $E_{\delta m_e}$ are negligible across observed scales, as expected

[Only the spatial gradient portions were computable from data]



$\rho_{i/e}$ \rightarrow ion/electron gyroradius k [km^{-1}] Gradient computation not reliable at scales smaller than MMS sep.
 $d_{i/e}$ \rightarrow ion/electron inertial length
 MMS sep. \rightarrow spacecraft separation

Hall and Electron Pressure Terms

Interplay between E_{Hall} and E_{Pe} important because E_{Pe} can provided a *non-ideal* electric field that can energise electrons and contribute to dissipation

E_{Hall} dominates over E_{Pe} , but E_{Pe} is stronger than expected from typical kinetic Alfvén wave predictions (**Fig. 5a**) [Boldyrev+ (2013) ApJ]
 → magnetic reconnection may lead to excess E_{Pe} [Stawarz+ (2019) ApJL]

Partial anti-alignment between E_{Hall} and E_{Pe} fluctuations at sub-ion scales (**Fig. 5b**) → degree of alignment linked to relative importance of ion and electron dynamics in supporting sub-ion scale currents

Linear & Nonlinear Terms

$$E_{MHD} = \boxed{-\delta u \times B_0} \boxed{-\delta u \times \delta b} \quad E_{Hall} = \boxed{\delta j \times B_0 / en} + \boxed{\delta j \times \delta b / en}$$

In E_{MHD} (**Fig. 6a**), nonlinear to linear term ratio increases at sub-ion scales → caused by decrease in alignment of δu and δb at sub-ion scales [consistent with Parashar+ (2018) PRL]

In E_{Hall} (**Fig. 6b**), ratio of nonlinear to linear terms is $\sim \delta b_{rms}/B_0$ at all scales

Results suggest a balance of linear and nonlinear timescales at both MHD and sub-ion scales that is set by $\delta b_{rms}/B_0$

Figure 5

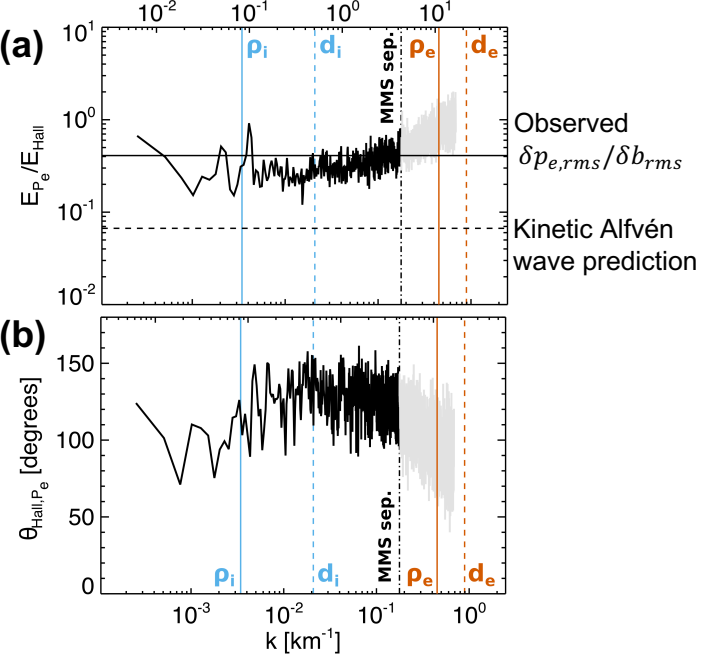


Figure 6

